

SEESAW MECHANISM AND ITS VARIOUS REALIZATIONS



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2. TYPE I , TYPE II AND DOUBLE SEESAW
3. LARGE MIXINGS FROM TYPE II SEESAW : TWO EXAMPLES
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 - A) 3×2
 - B) SINGULAR

SEESAW MECHANISM FOR SMALL NEUTRINO MASSES:

STD MODEL + ν_R WITH A
LARGE MAJORANA
MASS ($B-L=2$)

$$\begin{matrix} \nu_L & \nu_R \\ \nu_L & \begin{pmatrix} 0 & M_D \\ M_D^T & M_R \end{pmatrix} \\ \nu_R & \end{matrix}$$



$$\Rightarrow \mathcal{M}_\nu \approx -M_D^T M_R^{-1} M_D$$

$$M_D \leq v_{wk} \ll M_R \Rightarrow m_\nu \ll m_{e,u,d}$$

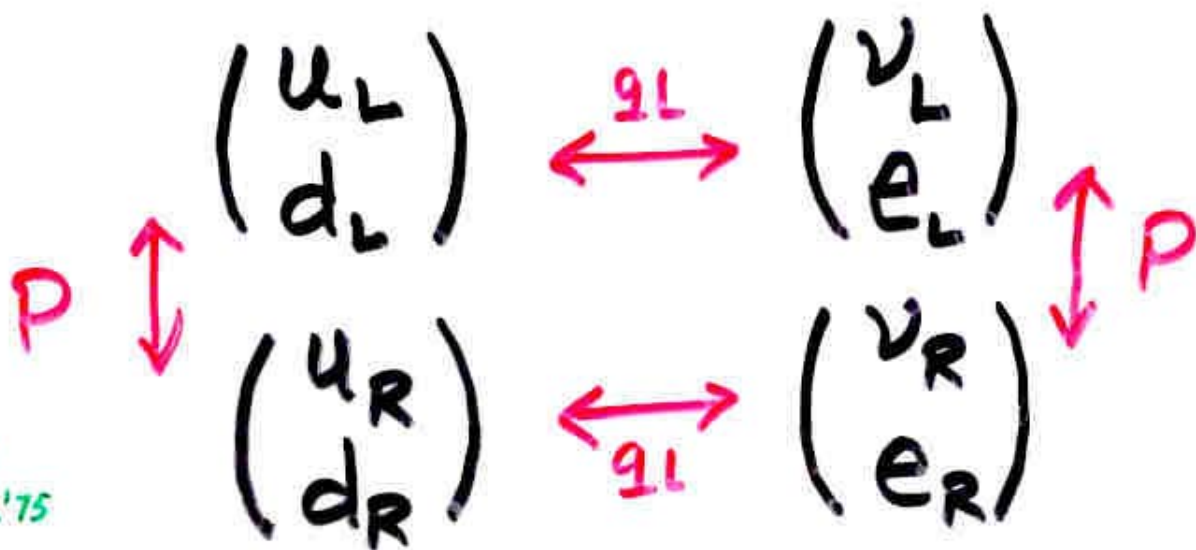
TYPE I SEESAW

GELL-MANN, RAMOND, SLANSKY; YANAGIDA; R.N.M., SENJANOVIĆ, GLASHOW
(1979)

PHYSICS OF THE SEESAW MECHANISM

———— X ————
 ADDING ν_R TO STD MODEL OPENS UP A WHOLE NEW LANDSCAPE OF PHYSICS !!

- RESTORES BOTH QUARK-LEPTON
- WEAK INT. BECOMES PARITY CONSERVING



PATI, SALAM;
 R.N.M. PATI '74
 SENJANOVIĆ, R.N.M. '75

- GAUGE SYM. $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$

$$Q = I_{3L} + I_{3R} + \frac{B-L}{2}$$

A. DAVIDSON
 R.N.M., MARSHAK 2/79

TWO ISSUES FOR LEFT-RIGHT SYM. MODELS:

_____ x _____

(i) HOW TO GET TO STD MODEL?
 $SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_X$



$SU(3)_c \times SU(2)_L \times U(1)_Y$

↓ $\langle \Phi(2, 2, 0) \rangle$

$SU(3)_c \times U(1)_{em} \neq 0$

(ii) SMALLNESS OF M_ν :

SINCE: $h_q \bar{Q}_L \Phi Q_R + h_\nu \bar{\Psi}_L \Phi \Psi_R$

$\Rightarrow m_{q, l, \nu}$ OF SAME ORDER!!

● M_ν AND SPONTANEOUS BREAKING OF PARITY:

$$SU(2)_L \times SU(2)_R \times U(1)_X$$

$$\begin{matrix} \nu_L & N_R \\ \left(\begin{array}{cc} 0 & 0 \\ 0 & f v_R \end{array} \right) \end{matrix}$$

$$v_R = \langle \Delta_R^0 \rangle$$

$$SU(2)_L \times U(1)_Y$$

$$\begin{matrix} \nu_L & N_R \\ \left(\begin{array}{cc} f v_L & m_{\nu D} \\ m_{\nu D} & f v_R \end{array} \right) \end{matrix}$$

$$v_{wk} \equiv \langle \Phi \rangle$$

$$U(1)_{em}$$

$$m_{\nu D} \approx h v_{wk} \ll f v_R$$

$$m_\nu = f \frac{v_{wk}^2}{2v_R} - M_{\nu D}^T \frac{f}{v_R} M_{\nu D}$$

PARITY SYM. \Rightarrow BOTH ν_L & ν_R

THREE IMPLICATIONS

- (i) ν_R MASS CONTROLLED BY B-L SYM. SCALE
- (ii) $m_\nu \rightarrow 0$ AS WEAK INT. BECOMES PURE V-A.
- (iii) MODIFICATION OF SEESAW (TYPE II)

PARITY SYMMETRY

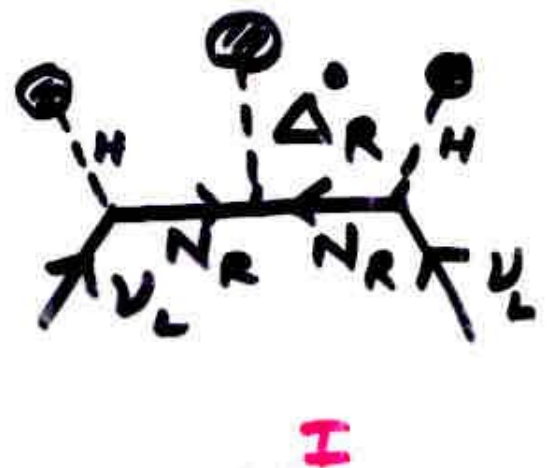
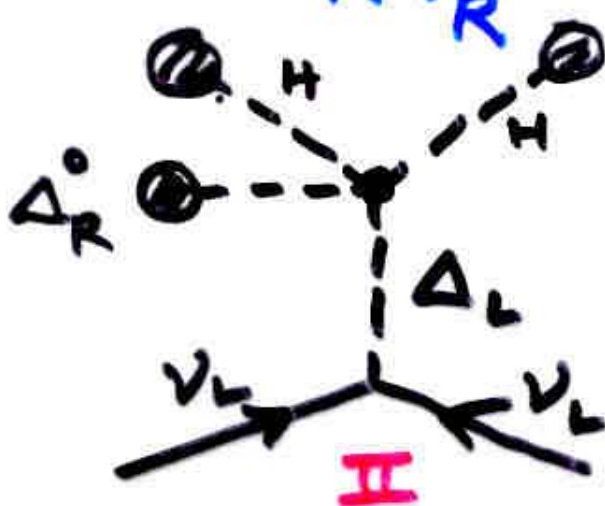
AND TYPE II SEESAW:

PARITY \Rightarrow $\nu_R \nu_R \Delta_R \leftrightarrow \nu_L \nu_L \Delta_L$
 OR $(\nu_R \chi_R)^2 \leftrightarrow (\nu_L \chi_L)^2$

TYPE II

LAZARIDIS, SHAFI; WETTERICH '80
 R.N.M., SENJANOVIC '80

$$M_\nu = f \frac{v_{wk}^2}{2v_R} - M_D^T (f v_R)^{-1} M_D$$



OTHER REALIZATIONS
OF SEESAW:

TEV SCALE DOUBLE SEESAW

(R. N. M., VALLE '86)

TYPE I OR II SEESAW + Δm_{ATM}^2

$$\Rightarrow M_R \sim 10^{15} \text{ GeV}$$

IS IT ALWAYS TRUE?

\Rightarrow BEYOND ν_R : $\nu_R \gg S$.
(AS IN E_6)

$$\Rightarrow \begin{pmatrix} \nu_L & \nu_R & S \\ 0 & M_{\nu D} & 0 \\ m_{\nu D}^T & 0 & M \\ 0 & M^T & \mu \end{pmatrix}$$

$$\Rightarrow M_\nu = -m_{\nu D} M^{-1} \mu M^T m_{\nu D}^T$$

$\Rightarrow m_{\nu D}, \mu \ll M \Rightarrow M \sim \text{few TeV}$
(EXPECT $\text{TeV } \tilde{Z}'$) POSSIBLE.

DOUBLE SEESAW TO TYPE III SEESAW:

BARR '03
ALBRIGHT, BARR '04

PARITY SYM:

$$\begin{pmatrix} 0 & m_{\nu D} & f\nu_L \\ m_{\nu D}^T & 0 & f\nu_R \\ f\nu_L & f\nu_R & M' \end{pmatrix}$$

$$\Rightarrow m_{\nu} = -m_{\nu D} M'^{-1} M' M'^{-1} m_{\nu D}^T$$

$$= \frac{(m_{\nu D} + m_{\nu D}^T) \nu_L}{M}$$

$$M \equiv f\nu_R \sim M' \sim 10^{15} \text{ GeV}$$

SEESAW IN THE QUARK SECTOR

$SU(2)_L \times U(1)_Y$ SINGLET VECTOR LIKE QUARKS, Q (IN LR)

$$\begin{matrix} q_L & Q_L \\ q_R & \begin{pmatrix} 0 & h v_R \\ h v_L & M \end{pmatrix} \\ Q_R & \end{matrix} \Rightarrow m_q \sim \frac{h^2 v_L v_R}{M}$$

LESS FINE TUNING OF YUKAWAS:

DAVIDSON, WALI '87
'88.

⋮

HOW USEFUL IS SEESAW FOR UNDERSTANDING NEUTRINO MASSES AND MIXINGS?

OF PARAMETERS: $18 + 3 = 21$

OF γ -OBSERVABLES: $3 + 3 + 3$
 m_i θ_{ij} δ_i

⇒ MODEL BUILDING: (CP CONSERVING)

- (i) $SO(10)$ (TYPE II) (1-PARAM)
- (ii) RADIATIVE CORRECTIONS (TYPE II) (3-)
- (iii) 3×2 SEESAW (7)
- (iv) $L_e - L_\mu - L_\tau$ AND SINGULAR SEESAW. (4-)

MINIMAL $SO(10)$ WITH $\Delta(B-L) = 2$ BREAKING:

(SUSY)

_____ X _____

BABU, R.N.M. '92
BAJC, SENJANOVIĆ, VISSANI '02
GOH, R.N.M., NG '03.

$SO(10)$: MATTER $\psi \subset \{16\}$

HIGGS: $\{210\}$ BREAKS $SO(10)$.

$\Delta = \{126\}$ $\Rightarrow SU(2)_R \times U(1)_B$

$H = \{10\}$ $\Rightarrow SU(2)_L \times U(1)_Y$

YUKAWA: $W = h \psi \psi H + f \psi \psi \bar{\Delta}$

- f UNIFIES ν 'S WITH q, Q . FLAVORS.
- 12 PARAMETERS (SOFT ~~SP~~)
ALL FIXED BY (Q, e, μ, τ) .

$$M_u = h v_u + f k_u$$

$$M_d = h v_d + f k_d$$

$$M_e = h v_d - 3f k_d$$

$$M_{\nu D} = h v_u - 3f k_u$$

$$M_{N_R} = f v_R.$$

$$\langle \Phi_{10} \rangle_{u,d} = v_{u,d} ; \langle \Phi_{126} \rangle_{u,d} = k_{u,d}.$$

TYPE II SEESAW:

$$M_\nu = f \frac{v_{wk}^2}{\lambda v_R} - M_{\nu D} \frac{1}{f v_R} M_{\nu D}$$

FOR $\lambda \ll 1$
 $v_R = M_U$; $M_\nu = m_0 f = c (M_d - M_e)$
(RELATION AT M_U).

MUST EXTRAPOLATE $m_{q,e}$ TO M_U !!

$$M_{d,b}^{(M_U)} \approx m_{b,\tau}^{(M_U)} \begin{pmatrix} \lambda^4 & \lambda^4 & \lambda^3 \\ \lambda^4 & \lambda^2 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix}$$

PHENOMENOLOGICAL FACT:

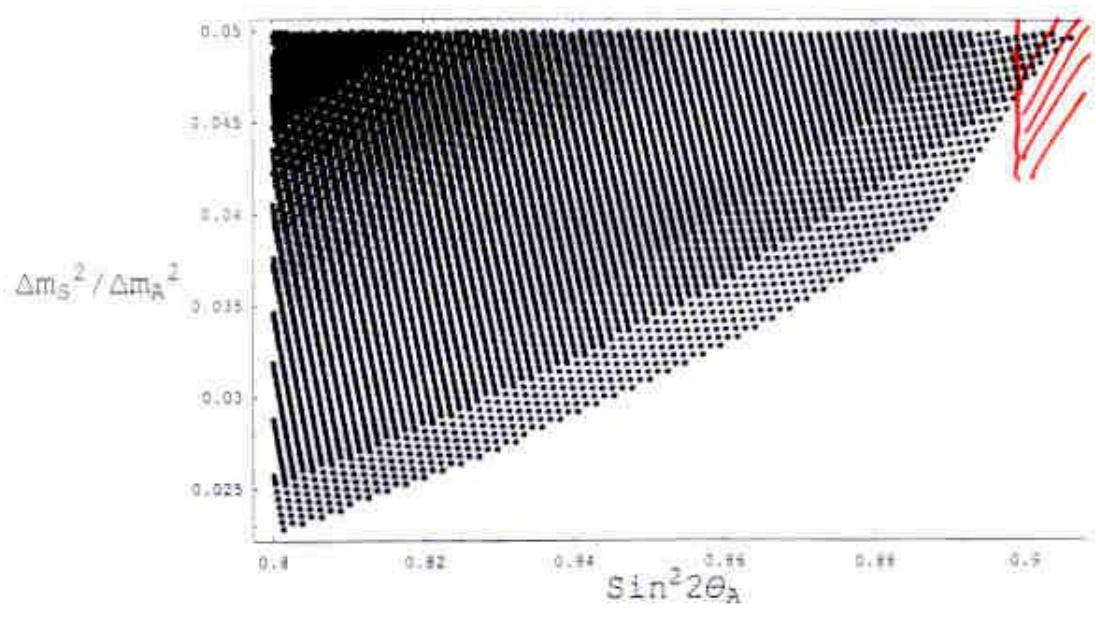
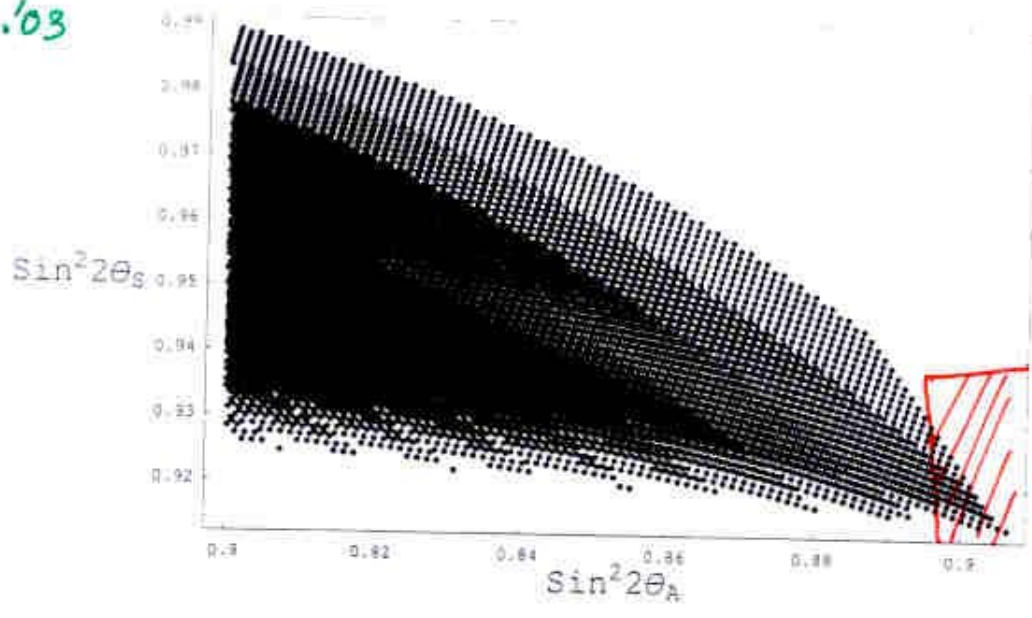
$$m_b^{(M_U)} \approx m_\tau^{(M_U)} + O(\lambda^2)$$

$$\Rightarrow M_\nu \approx c (M_d - M_\ell)$$

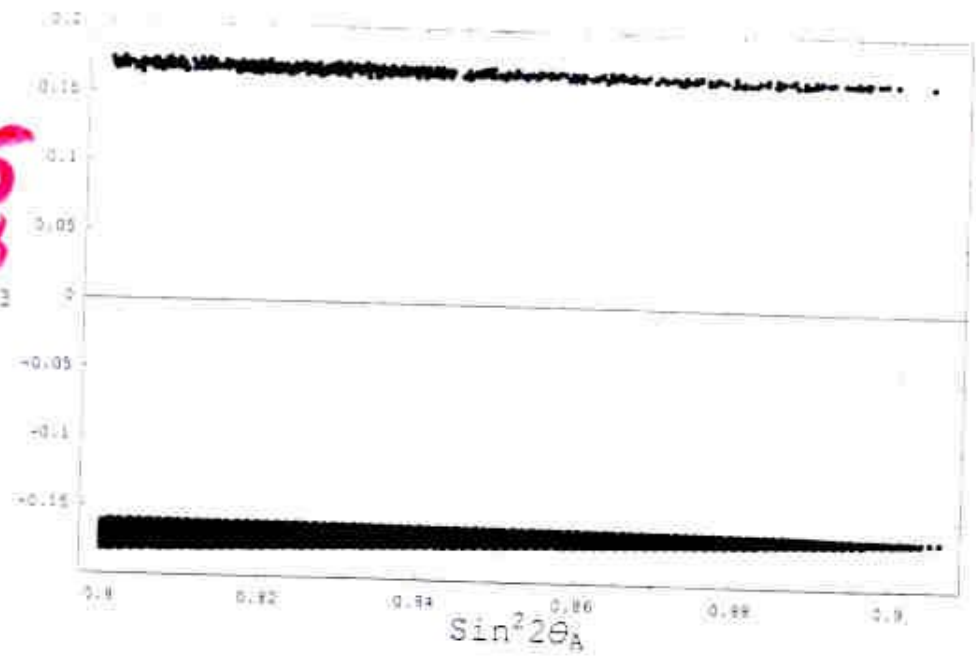
$$\approx m_0 \begin{pmatrix} \lambda^4 & \lambda^4 & \lambda^3 \\ \lambda^4 & \lambda^2 & \lambda^2 \\ \lambda^3 & \lambda^2 & \lambda^2 \end{pmatrix}$$

- \Rightarrow
- θ_{12} LARGE ; • θ_{23} LARGE
 - $\theta_{13} \approx \lambda$; • $\sqrt{\frac{\Delta m_{21}^2}{\Delta m_{31}^2}} \approx \lambda$.

$\tan \beta = 10$



$U_{e3} \approx 0.16$
 $U_{e3} \approx 0.18$



TESTABLE
IN
LBL
EXPTS.

TYPE I SEESAW FOR
THIS MODEL DOES NOT WORK !!

FUKUYAMA, OKADA '02
MATSUDA, KOIDE, FUKUYAMA, NISHIURA '02
DUTTA, MIMURA, R.N.M. '04

LARGE MIXING FOR QUASI-DEG. ν 'S:

X

BALAZI, DIGHE, R.N.M., PARIDA '00
ANTUSCH, KERSTEN, LINDNER, RATZ '02
R.N.M., PARIDA, RAJASEKARAN '03

Q-L UNIFICATION

→ AT SEESAW SCALE

$$\theta_{\text{quark}} \approx \theta_{\text{lepton}}$$

THIS CAN HAPPEN FOR THE TYPE II SEESAW CASE:

$$M_\nu = f \nu_L - M_{\nu D} \frac{1}{f \nu_R} M_{\nu D}^T$$

SUPPOSE $f \approx 1 + O(\epsilon)$

→

$$m_1 \approx m_2 \approx m_3$$

AND

$$\theta_q \approx \theta_l$$

AT SEESAW SCALE

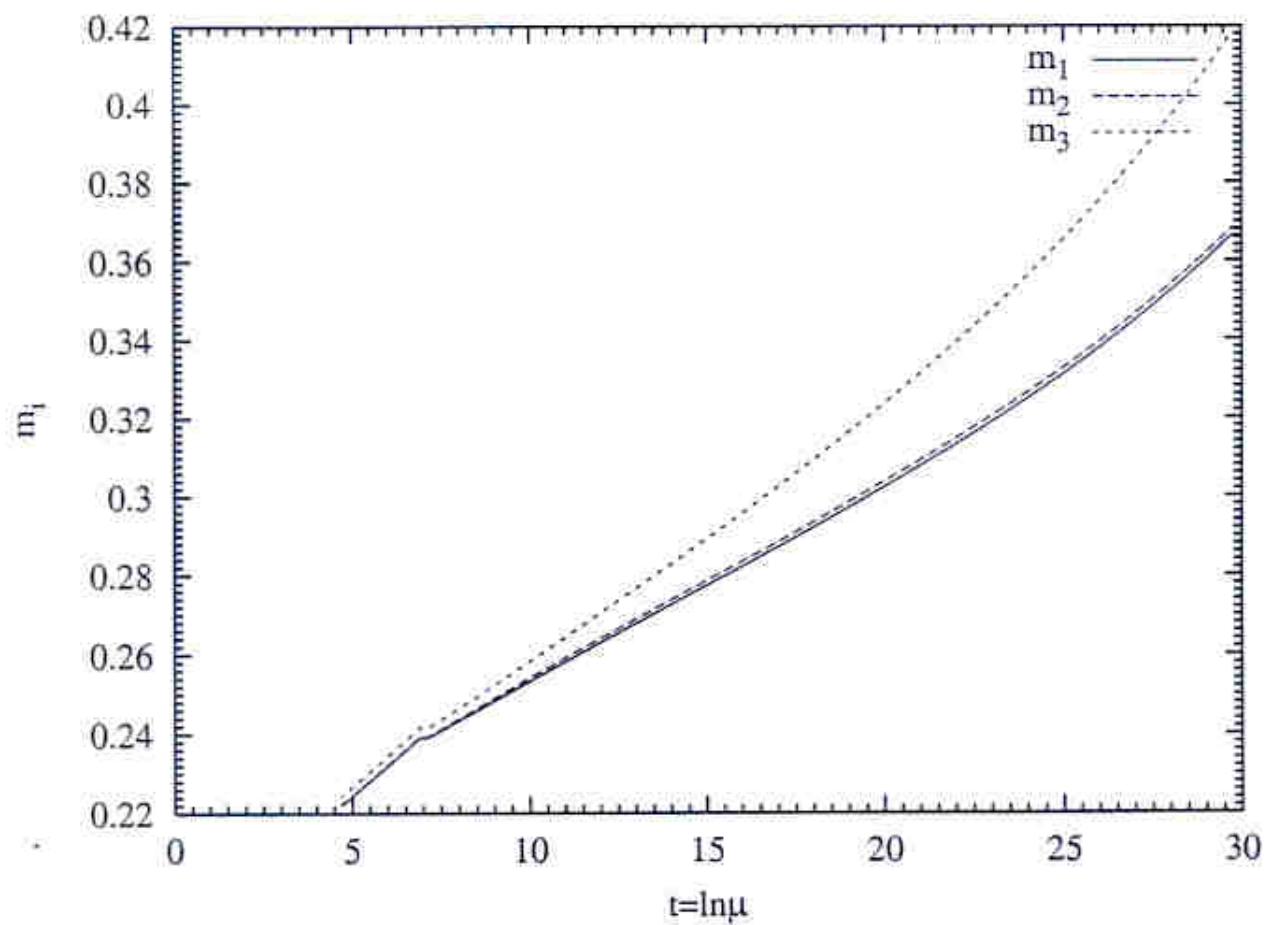
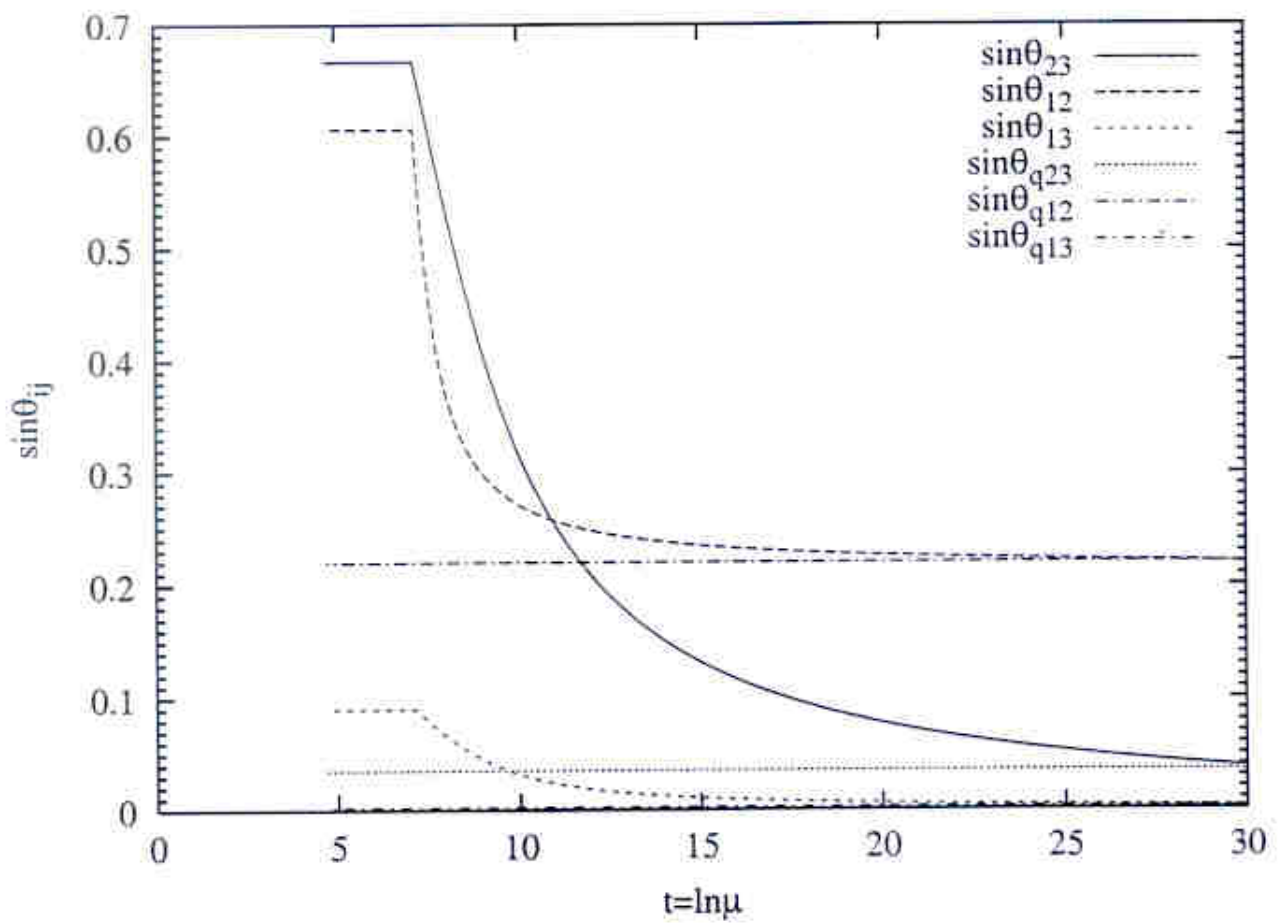
AT WEAK SCALE:

$$M_{\nu}(M_Z) = \begin{pmatrix} 1+\delta & \\ & \end{pmatrix} M_{\nu}(M_U) \begin{pmatrix} \\ & 1+\delta \end{pmatrix}$$

$$\frac{d\theta_{ij}}{dt} \propto \frac{m_i + m_j}{m_i - m_j} \theta_{ij}$$

CASAS, IBARRA,
ESPINOSA, NAVARO

LARGE SLOPE FOR $m_2 \approx m_3$



PARIDA, RAJASEKARAN, R.N.M.
'03

• SEESAW WITH 2 RH NEUTRINOS: (3x2 SEESAW)

• LOCAL B-L REQUIRES 3 N_R 'S FOR ANOMALY CANCELLATION

2 N_R ALTERNATIVES:

(A) $M_{R3} \gg M_{R1}, M_{R2}$

SMIRNOV '93
GRIMUS & LAVOURA '00
KING '00
KING, ROSS '01

$$\Rightarrow M_\nu = - M_{\nu D}^T M_R^{-1} M_{\nu D}$$

3x2 2x2 2x3

(B) ONLY 2 N_R 'S: KUCHIMANCHI, R.N.M. '02
RABY '03

GAUGE GROUP: $G_{STD} \times SU(2)_H$

GLOBAL WITTEN ANOMALY CANCELLATION
REQUIRES 2 RH ν_R 'S.

ADVANTAGES AND PREDICTIONS

LESS PARAMETERS:

3×3	CP 21	CP 12
3×2	14	8 7 ($L_e - L_\mu$)

PREDICTIONS:

• INVERTED HIERARCHY

$$\Delta m_{23}^2 < 0$$

$$\theta_{13} = 0$$

INTERESTING COSMOLOGY:

• FRAMPTON, GLASHOW, YANAGIDA '02

• RAIDAL, STRUMIA '02; ENDOH et. al. '02; AKHMEZOV, FRIGERIO, SMIRNOV '03

PHENOMENOLOGY:

CHANKOWSKI, ELLIS, POKORSKI, RAIDAL, TURCZYNSKI '04

DUTTA, R.N.M. '03

IBARRA, ROSS '04

$$\textcircled{c} M_{R3} \lll M_{R1}, M_{R2}.$$

NATURAL IF TH. $L_e - L_\mu - L_\tau$ INV.
(R.N.M. '01)

$$\Rightarrow M_{NR} = \begin{pmatrix} 0 & M_1 & M_2 \\ M_1 & 0 & 0 \\ M_2 & 0 & 0 \end{pmatrix}$$

$$\text{Det. } M_{NR} = 0.$$

SINGULAR:
SEESAW $M_\nu = -M_{\nu D} M_{NR}^{-1} M_{\nu D}^T$

HOWEVER, IF $\text{Det } M_{\nu D} = 0$

\Rightarrow SEESAW WORKS.

LIGHT ν - SPECTRUM: (LOWEST ORDER)

2 MASSIVE

1 MASSLESS

+ 1 ν_R MASSLESS.

\Rightarrow 3+1 SCENARIO FOR LSND!!

IF MINI-BOONE CONFIRMS LSND \rightarrow USEFUL MODEL.

OTHER EXAMPLES:

2 + 2 EXAMPLE

CHOUBEY, BRAHMACHARI, RNM
'02

3 + 2 EXAMPLE

GOLDMAN, MCKELLAR,
STEPHENSON, GARBUTT
'04
McDONALD, MCKELLER '09