

SEE-SAW

and

GRAND UNIFICATION

Paris, June 2004

SEE SAW MOST COMMONLY DEFINED AS

(IN THE SM)

Gell-Mann, Ramond,
Sleazy

add $\nu_R \dots$

Yonekida

Mohapatra, G.S.

Glashow (79)

$$\mathcal{L}_\nu(\nu) = \gamma_0 \bar{\nu}_L \phi \nu_R + M_R \nu_R^T C \nu_R + \text{h.c.}$$

$$\begin{matrix} \nu_L \\ \nu_R \end{matrix} \begin{pmatrix} 0 & m_D \\ m_D & M_R \end{pmatrix}$$

$$m_D = \gamma_0 \langle \phi \rangle \lesssim M_W$$

a mess, unless $M_R \gg M_W$

NATURAL : M_R a gauge invariant quantity

(minimal fine tuning)

$$m_\nu^I = m_D^T M_R^{-1} m_D \ll M_W, m_f \text{ (type I see-saw)}$$

NEUTRINOS EXPECTED TO BE

LIGHT

$$10^{10} \text{ GeV} < M_R < 10^{15} \text{ GeV}$$

ALTERNATIVELY

Add $SU(2)_L$, $B-L=2$ triplet Δ

$$\Delta L_Y = Y_0 \ell_L^T C \ell_L \Delta \quad \ell = \begin{pmatrix} \nu \\ e \end{pmatrix}$$

$$\Delta V = M \bar{\Phi}^T \Delta \Phi + M_\Delta \Delta^+ \Delta + \dots \text{O.T.}$$

Happ, Wetterich '80

$$\langle \Delta \rangle \approx \frac{M_W^2 M}{M_\Delta^2} \approx \frac{M_W^2}{M_\Delta}$$

Logarithms,
in ϕ_i , ψ_i

Mohapatra, G.S. '81

FOR $M \sim M_\Delta \gg M_W$ (same principle as before)



$$m_\nu^{\text{II}} = Y_0 \frac{M_W^2}{M_\Delta} \ll M_W, m_f \quad \text{NATURAL}$$

Type II see-saw

~~SM: no difference; an effective operator~~

~~$$\frac{\ell_L^T C \Phi^+ \Phi^+ \ell_L}{M_\nu}$$~~

~~M_ν~~

~~Need $Y_0, M_R, Y_\nu \dots$~~

~~Weinberg~~

~~$$M_\nu \ll M_{\text{pl}} ??$$~~

SM : NO DIFFERENCE

Needs a theory of mass scales;

(restriction on) Y_D, Y_U, \dots

e.g. $Y_D \propto Y_U \propto M_R \Rightarrow \text{same } (I = \bar{U})$

SM - effective operator

$$\Delta L_y = \frac{(\bar{l}_i^T c \Phi^*) (\Phi^+ \sigma_2 l_j)}{M_\nu} c_{ij}$$

Weinberg '79

$$m_\nu < 1 \text{ eV}$$

$$M_\nu : 10^{10} - 10^{15} \text{ GeV}$$

$$\begin{array}{l} \rightarrow M_\nu \ll M_{\text{pl}} \quad (M_{\text{pl}} \Rightarrow m_\nu \approx 10^{-5} - 10^{-6} \text{ eV}) \\ \text{WHY?} \quad \Downarrow \end{array}$$

GRAND UNIFICATION : a theory
of mass scales (large)

+ restriction on Yukawa

How to incorporate see-saw
in GUTS ?



WHAT IS 'THE' GUT ?

Minimal consistent theory
NO D-T... (CORRECT LOW ENERGY THEORY)

~~CREATED~~ DARK

Old days : $m_\nu = 0$ prejudice !



minimal SU(5) + supersymmetry



unification
(hierarchy)

Mostly SUSY GUTS ;

(ordinary at the end)

'PURE' GUTS (no additional symmetries)

SYMMETRY REASONING

Minimal SH \Rightarrow

accidental global B-L

↑ ANOMALY FREE

broken by M_R (H)

- MAYBE ACCIDENT \Rightarrow

 $SU(5)$ DIRECTION

- GAUGE B-L $\Rightarrow (B-L)^3$ anomaly

 $\Rightarrow 3 U_R$ (simplest)

$$SU(5) \times U(1) \cong \begin{matrix} SO(10) & \text{DIRECTION} \\ U(1) & \end{matrix} \quad (L-R; \text{Pati-Salam})$$

$$SU(2)_L \times SU(2)_R \times SU(4)_C$$

$$U(1)$$

$$SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

$$Q = T_{3L} + T_{3R} + \frac{B-L}{2}$$

Molodtsov, G.S.

GRAND UNIFICATION

• SU(5) THEORY

give up gauging B-L

↑
accidental global symmetry

neutrino mass:

(i) $\exists \nu_R$ - singlets of SU(5)

Minimal fine-tuning:

M_R gauge invariant $\Rightarrow M_R \gg M_{GUT} \approx 10^{16}$ GeV
(susy)

$$\Rightarrow m_\nu \ll \frac{M_W^2}{M_{GUT}} \approx 10^{-3} \text{ eV}$$

↑ SMALL

(ii) add $\Delta \leq 15$ (symmetric 5×5)

same principle $M_\Delta \gg M_{GUT}$

NOT RIGOROUS - but in the spirit
of see-saw

such
SU(5) not a good theory of fermion masses
minimal $\Rightarrow m_d = m_e$ at M_{GUT}

$m_b \approx m_\tau$ at M_{GUT}

$m_s = 1/3 m_\mu$
 $m_d = 3 m_e$ } large deviations

(45) - fermi-jarlskog

Needs new Higgs or /and

non-renormalizable interactions ($1/M_{pl}$) Barbieri et al



NOT PREDICTIVE

POSSIBILITY: add extra (flavor) symmetries

• $d = 5$ proton decay ?

ruled out : Murayama, Pierce ('01)

no : Bqic, Fileviez Pires, G.S. ('02)

● $SO(10)$: MINIMAL GUT OF MATTER
AND INTERACTIONS

ideal framework for \mathcal{D} masses
and mixings (other fermions)

- a family of fermions $= 16$ (spinor)

$$16 = \begin{pmatrix} f \\ f^c \end{pmatrix}_L \quad (f^c)_L \equiv c \bar{f}_R^T$$

- L-R a gauge symmetry : G

- supersymmetry R-parity (matter parity)
a gauge symmetry
Hohapatra '86

$$R: p \rightarrow p, \bar{p} \rightarrow -\bar{p} \quad R = (-1)^{2J} M \equiv \begin{matrix} 2J \\ 3(B-L) \\ (-1) \end{matrix}$$

$$M: f \rightarrow -f, f^c \rightarrow -f^c, \Phi \rightarrow \bar{\Phi}$$

$$16 \rightarrow -16, 10 \rightarrow 10$$

$$\hookrightarrow \subseteq Z_4 \text{ center of } SO(10) \\ (16 \rightarrow i16, 10 \rightarrow -10)$$

Yukawa sector in $SO(10)$

$$16_F \times 16_F = 10_H + 120_H + 126_H$$

$$16_F \subset \Gamma_i: 16_F \ 10_i \quad \text{symmetric}$$

$$16_F \subset \Gamma_i: \Gamma_j \Gamma_k 16_F \ 120_{[ij\bar{k}]} \quad \text{anti}$$

$$16_F \subset \Gamma_i \Gamma_j \Gamma_k \Gamma_l 16_F \ \overline{126}_{[ij\bar{k}lm]} \quad \text{symmetric}$$

which is correct?

What is the minimal Higgs?

Natural: ONLY $10_H \leftarrow$ generates M_W
and m_f (SM)

see-saw \Rightarrow more Higgs: $\overline{126}$

STRONGER: $\overline{126}$ (direct or effective)

needed for charged fermion masses
(10 NOT ENOUGH)

Pati - Salam decomposition : $SU(2)_L \times SU(2)_R \times SU(4)_C$

$$\begin{pmatrix} u & u & u & \nu \\ d & d & d & e \end{pmatrix}_{L,R}$$

$$16 = (2, 1, 4) + (1, 2, \bar{4})$$

e, l e^c, l^c

useful language!

$$10 = (2, 2, 1) + (1, 1, 6)$$

↑
LIGHT HIGGS

$$\bar{126} = \underline{(3, 1, 10)} + \underline{(1, 3, \bar{10})} + \underline{(2, 2, 15)} + (1, 1, 6)$$

↑ ↑ ↑
see saw II? mass t, ν_R LIGHT HIGGS?

NOTICE: $\langle 10 \rangle = \langle 2, 2, 1 \rangle \sim$ PS singlet



✓
 $\theta_2 = 0$

$m_2 = m_l$ at M_{GUT}
(even if more 10's)

Fails badly for 2nd and 1st generation

⇒ $\bar{126}$ A MUST !!

$\langle 2, 2, 15 \rangle$:

$3 m_2 = -m_e$
(2nd gen.)

(tree level; loop;
renormalizable; non-renorm.?)

Gemzi - Glesner decomposition $SU(5)$

$$16 = 10 + \bar{5} + 1 \rightsquigarrow \nu^c$$

$$10 = 5 + \bar{5} \leftarrow \text{light Higgs}$$

$$\overline{126} = 1 + 5 + 15 + \overline{45} + 50$$

\uparrow mass for ν^c \uparrow types? \uparrow light Higgs

(Gemzi - Jarkslög)

• effective $\overline{126}$

$$\frac{1}{M_{pe}} \overline{16}_H \text{ C } \Gamma \Gamma \Gamma \Gamma \Gamma \overline{16}_H \sim \overline{126}_H$$

NON-RENORMALIZABLE
COUPLINGS

MINIMAL SO(10) ?

MINIMAL GUT ?

● SMALL REPRESENTATIONS

$(\psi, \bar{\psi})$ 16_H (+ $\bar{16}_H$ in SUSY) } M_{GUT}
 (A) 45_H - adjoint } $M_G?$

(H) 10_H - low scale (M_w)

small # fields \Rightarrow • AF above M_{GUT}
 • high precision(?)

SUSY

● COMPLICATED HIGGS SECTOR

$$\begin{aligned}
 W_H = & m_A^2 A^2 + \frac{(A^2)^2 + A^4}{M_{pe}} \\
 & + m_\psi \psi \bar{\psi} + \frac{(\psi \bar{\psi})^2 + (\psi \Gamma \Gamma \bar{\psi})^2 + (\psi (\Gamma \Gamma \Gamma) \bar{\psi})^2}{M_{pe}} \\
 & + \frac{(\psi \Gamma \psi)^2 + (\psi \Gamma^5 \bar{\psi})^2 + \psi \rightarrow \bar{\psi}}{M_{pe}} \\
 & + \frac{\psi \Gamma \Gamma \psi A + \bar{\psi} \Gamma \Gamma \bar{\psi} A}{M_{pe}} + (\psi \Gamma^4 \psi A^2 + \psi \rightarrow \bar{\psi}) \\
 & + m_H H^2 + \psi \Gamma \psi H + \bar{\psi} \Gamma \bar{\psi} H + \frac{\psi \Gamma^3 \psi A H}{M_{pe}} + \psi \rightarrow \bar{\psi}
 \end{aligned}$$

MORE PROBLEMATIC : PROTON DECAY ISSUE

$$\langle 16_H \rangle \neq 0 \neq \langle \overline{16}_H \rangle$$

BREAKS MATTER
(R) PARITY at M_{GUT}

DISASTER

$$16_F \Gamma 16_H 10_H \Rightarrow M_{GUT} L \overline{H}$$

$$\frac{(16_F \Gamma 16_F)(16_F \Gamma 16_H)}{M_{pe}} \rightarrow \frac{M_{GUT}}{M_{pe}} (U^c D^c D^c + Q L D^c)$$

10^{-3}

$d=4$ p decay
 $(m_{\tilde{d}^c} \approx TeV)$

\sim NEED $\leq 10^{-12}$



Ad-hoc new 'matter' parity

$$16_H \rightarrow -16_H, \quad \overline{16}_H \rightarrow -\overline{16}_H$$

UGLY !

Matter parity - gauge symmetry in $SO(10)$

• Yukawa sector messy

$$\begin{aligned}
 W_Y = & 16_F \Gamma 16_F H_{(10)} \leftarrow \text{top (b) quark mass} \\
 & + \frac{16_F \Gamma 16_F \Psi \Gamma \Psi}{M_{pe}} + \frac{16_F \Gamma 16_F \bar{\Psi} \Gamma \bar{\Psi}}{M_{pe}} \quad (10) \\
 & + \frac{16_F \Gamma \Gamma \Gamma 16_F A H}{M_{pe}} \quad (120) \leftarrow \text{mass corrections} \\
 & \quad \quad \quad (2, 2, 15) - GJ \\
 & + \frac{16_F \Gamma^5 16_F \bar{\Psi} \Gamma^5 \bar{\Psi}}{M_{pe}} \quad (\bar{126})
 \end{aligned}$$

$(2, 2, 15)$ \nearrow
 $(\text{mass corrections } - GJ)$ \nwarrow ν^c mass
 - automatically small (keeps b- τ unification)

5 SETS OF YUKAWAS

- Add flavor symmetries (textures) Abelglt
Babu, Barr
Raby ...

- $M_R = \frac{M_{GUT}^2}{M_{pe}} \approx 10^{13} \text{ GeV} \leftarrow \text{nice}$
- type II suppressed ($1/M_{pe}$)

BUT PROTON DECAY STRIKES AGAIN

$$\frac{1}{M_{\text{pl}}^2} (16_F \Gamma 16_F) (16_F \Gamma 16_F)$$



$$\frac{QQQL}{M_{\text{pl}}}$$

TOO FAST

$$\tau_p = 10^{23} \text{ yr}$$

Needs a large suppression:

$$C_{\text{pe}} \leq 10^{-6} \quad (\text{why others } O(1)?)$$

Could it be that $1/M_{\text{pl}}^2$
suppressed ???

... Maybe a renormalizable
theory only ??? !!!

elementary 126_H : GRS
Gleadow

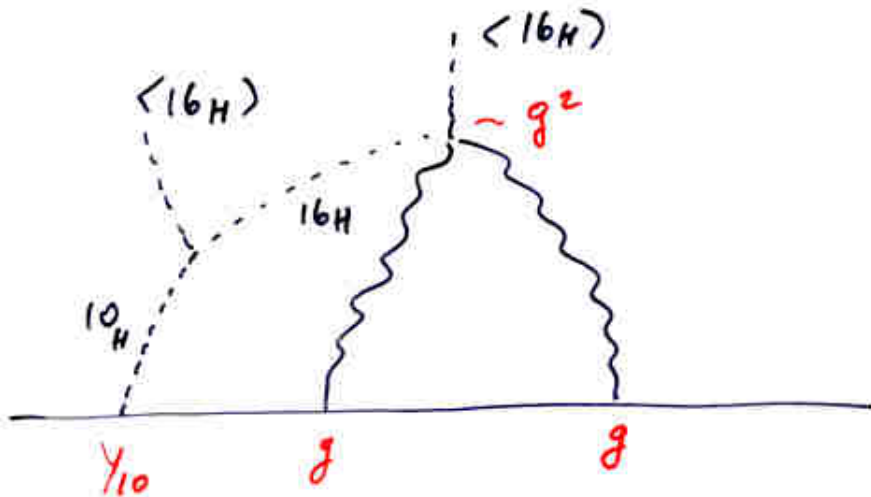
• (ILLUSTRATIVE) FAILURE

Witten '80

ORDINARY SO(10)

10_H, 16_H, 45_H

126 - EFFECTIVE TWO-LOOP



$M_{SUT} \sim 16_H \Gamma 16_H 10_H$

$$M_{VR} = Y_{10} \left(\frac{\alpha}{4\pi} \right)^2 M_{SUT}$$

$Y_{up} \leftarrow$ Pati-Salam $(2, 2, 1)$

10_H: $Y_0 = Y_{up}$, $M_{VR} \propto Y_{up}$

$m_\nu \propto M_{up}$

$m_e = m_{down}$

$\Rightarrow \theta_{e\mu\tau} = \theta_{\nu\mu\tau}$

$(= 0)$

MORE 10_H $\Rightarrow \theta_e \neq 0, \theta_{e\mu\tau} \neq 0$

(arbitrary) $M_{VR} \neq M_{up}$

$m_e = m_d \left(1 + \left(\frac{\alpha}{4\pi} \right)^2 \right)$

Message?

- SUSY LOVERS :

$$16_H 16_H 10_H \quad M_{GUT} \rightarrow M_{3/2} \quad (10 FT)$$

$$\Rightarrow m_{\nu_R} \ll M_{3/2} \quad \text{too small}$$

- FERMION (CHARGED) MASSES

$$\left(1 + \left(\frac{\alpha}{4\pi}\right)^2\right) m_d = m_e \quad \text{at } M_{GUT}$$

cannot be corrected in Witten

- only Y_{10} flavor structure
(motly)

126 - a must for charged fermions too

~~(a) (b) (c) (d) (e) (f) (g) (h) (i) (j) (k) (l) (m) (n) (o) (p) (q) (r) (s) (t) (u) (v) (w) (x) (y) (z)~~
~~QED~~

BIG IS BETTER

RENORMALIZABLE INT.

Clodas, Kuo, Nagayawa '82
Aulakh, Mohapatra '82

$$126_H + \overline{126}_H$$

$$210_H \text{ (4 index antisym.)}$$

$$10_H$$

Aulakh, Bajc, Heltzer,
Visseri, G.S. (2003)

(OR $54_H + 45_H$)

$$W_H = m_{210} (210)^2 + M_{126} 126 \cdot \overline{126}$$

$$+ \lambda (210)^3 + \eta 210 \cdot 126 \cdot \overline{126}$$

$$+ \kappa 210 \cdot 126 \cdot 10 + \overline{\kappa} 210 \cdot \overline{126} \cdot 10$$

$$W_Y = \gamma_{10} 16_F \Gamma 16_F 10_H + \gamma_{126} 16_F \Gamma^5 16_F \overline{126}_H$$

ONLY TWO SETS OF SYMMETRIC

(generation space) YUKAWAS



$$3 + 6 \cdot 2 = 15 \text{ real parameters (only)}$$

Pati-Salam decomposition

$$i = 1, \dots, 10$$

- 10_H : $10_a (2, 2, 1)$
 $10_\alpha (1, 1, \alpha)$

$$a = 1, \dots, 4 \quad SU(4) = SU(2)_c \times SU(2)_f$$

$$\alpha = 5, \dots, 10 \quad SU(6) = SU(4)_c$$

- $2 \cdot 10_H$:

$$[abcd] = \underline{(1, 1, 1)}$$

$$[abc]\alpha = (2, 2, 6)$$

$$[ab](\alpha\beta) = \underline{(1, 3, 15)} + (3, 1, 15)$$

$$a[\alpha\beta\gamma\delta] = \underline{(2, 2, 10)} + \underline{(2, 2, \bar{10})}$$

$$[\alpha\beta\gamma\delta] = \underline{(1, 1, 15)}$$

— = GUT vev fields

- $126_H, \overline{126}_H$:

$$[abcd]\alpha = (1, 1, 6)$$

$$[abc](\alpha\beta) = \underline{(2, 2, 15)}$$

$$[ab](\alpha\beta\gamma) = \underline{(1, 3, 10)} + (3, 1, 10)$$

↑
type II

uv = light Higgs candidates

Pati-Salam decomposition

$$i = 1, \dots, 10$$

- 10_H : $10_a (2, 2, 1)$
 $10_\alpha (1, 1, \alpha)$

$$\alpha = 1, \dots, 4 \quad SU(4) = SU(2)_c \times SU(2)_\alpha$$

$$\alpha = 5, \dots, 10 \quad SU(6) = SU(4)_c$$

- $2 \cdot 10_H$:

$$[abcd] = \underline{(1, 1, 1)}$$

$$[abc]\alpha = (2, 2, 6)$$

$$[ab](\alpha\beta) = \underline{(1, 3, 15)} + (3, 1, 15)$$

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$$[\alpha\beta\gamma\delta] = \underline{(1, 1, 15)}$$

— = GUT vev fields

- $126_H, \overline{126}_H$:

$$[abcd]\alpha = (1, 1, 6)$$

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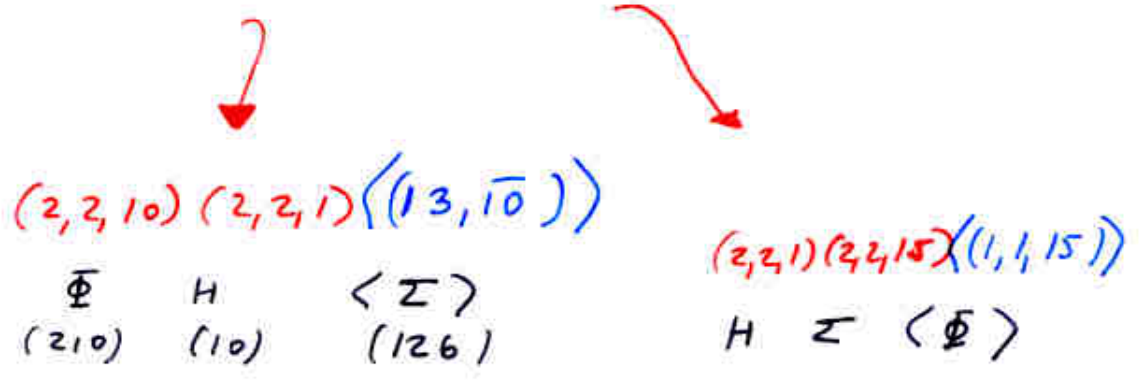
↑
type II

uv = light Higgs candidates

LIGHT HIGGS ?

$$\begin{aligned} (\Phi) &\equiv 210, \\ \Sigma, \bar{\Sigma} &\equiv 126, \bar{126} \\ H &\equiv 10 \end{aligned}$$

$$\Delta W_H = \alpha \underbrace{\Phi}_{210} \underbrace{\Sigma}_{126} \underbrace{H}_{10} + \bar{\alpha} \underbrace{\Phi}_{210} \underbrace{\bar{\Sigma}}_{126} \underbrace{H}_{10}$$



⇓

LIGHT Higgs = (2,2,1) + (2,2,15) + ...

$H(10)$ \uparrow $\bar{\Sigma}(126)$

⇓

$\bar{\Sigma}(126_H)$ at the same time

gives a mass to $\nu_R: \langle (13,10) \rangle$

and charged fermions: $\langle (2,2,15) \rangle$

Lasandes, Shafi, Wetterich '81

↑

Geri-Jarlskog

Babu, Mohapatra '92

See - saw: I or II (or both)?

(NEQUIVALENT

Bajc, Visseri, G.S.
'2006

I : $m_\nu^{\text{II}} \propto M_e - M_d$

$$\begin{pmatrix} m_\mu & 0 \\ 0 & m_\tau \end{pmatrix}$$

2-3 gen.

θ_e small $\Rightarrow m_\nu^{\text{II}} \propto \begin{pmatrix} m_\mu - m_s & \epsilon \\ \epsilon & m_\tau - m_b \end{pmatrix}$

Large $\theta_{atm} \Leftrightarrow m_b = m_\tau$

Bajc, Visseri, G.S.
PRL '2005

3 generations: numerical studies

large 1-3 mixing: $|U_{13}| \approx .15$

Gole, Holroyd, Ng
'2007

R-PARITY

23
Aulakh, Heltz, G.S.

better $M = (-1)^{3(B-L)} \equiv (-1)^{2S} R$

+ Rasin

• $M \overline{126} = \overline{126}$

• $\langle \tilde{\nu}_R \rangle = 0$ (in 16_F) 

$M \tilde{\nu}_R = M_R \gg m_{3/2}$ - does not
destabilized

• $\langle \tilde{\nu}_L \rangle = ?$

Aulakh, Heltz (82)

IN MSSM \Rightarrow MAJORON COUPLED

TO Z - ruled out!

integrate out N (νc) superfield

$$W_{\text{eff}} = Y_0^2 \frac{(LH)^2}{m_N} \Rightarrow V = m_{3/2}^2 Y_0^2 \frac{(\tilde{L}H)^2}{m_N}$$

$$\langle \tilde{\nu}_L \rangle = \langle \tilde{L} \rangle \neq 0 \Rightarrow m_{\text{Im } \tilde{\nu}_L}^2 = Y_0^2 \frac{m_{3/2}^2 m_W^2}{m_N} \ll m_W$$



Z decay \Rightarrow ruled out

R-parity exact

STABLE LSP (DARK
MATTER)

OBJECTIONS

- NOT asymptotically free

$$\Lambda_F \lesssim 10 M_{GUT} \text{ (strong coupling)}$$

- WHY NO $1/M_{Pl}$ TERMS ($1/\Lambda_F$)?



Hint: $c_{pe} \frac{QQQL}{M_{Pl}} \Rightarrow c_{pe} \leq 10^{-6}$

$$c_F \frac{QQQL}{\Lambda_F} \Rightarrow c_F \leq 10^{-8}$$

All small? OR hierarchy
large to small?

- computations involved (C-G),
but under control

Bajc, Melfo, Visseri, G.S.

Fukuyama, Iliakovac,
Kikuchi, Meljanac, Okada

← Aulakh, Girdhar

threshold effects
small?

(2004)

CONCLUSIONS

?

QUEST GOES
ON

LEPTO GENESIS

Rather involved; not clear who is
the culprit

most natural:

$\nu_R (\tilde{\nu}_R)$ decay κ

but Δ could naturally
dominate the process (Π type)

needs detailed study

(planned)

Hambye, G.S. '03

Antusch, King '04

Large M_R a problem in SUGRA

(gravitino abundance)

PROTON DECAY

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SUSY: • $d=4$ absent with $\overline{126}$
and 16 if more symmetry postulated

• $d=6$ somewhat slow

$$M_{GUT} \approx 10^{16} \text{ GeV} \Rightarrow \tau_p \approx 10^{36} \text{ yr}$$

• $d=5$ somewhat fast ($\overline{126}$)

needs a detailed study

Mohapatra et al
(2003)

Aulekh, Bajc, Heltzer,
Visseri, G.S.
(in progress)

ORDINARY:

$d=6$ possibly observable

Fukuyama et al
(04)

(fine-tuning)
8/6/2009

- $d=5$ proton decay:
uncertainties

$$16_H + \overline{16}_H, \quad 45_H$$

$$45 = \underbrace{24}_{SU(4)} + \underbrace{1}_{SU(5)} + \underbrace{10}_{SU(5)} + \overline{10}_{SU(5)}$$

$$24 = p_c + 3w + \text{"goldstones"}$$

NO CUBIC TERM $(24)^3 \Rightarrow$

$$\text{FROM } (24)^4: \quad m_3 = 4 m_8$$

$$\downarrow$$

$$M_{GUT} = M_{GUT}^0 \left(\frac{M_{GUT}^0}{2 m_8} \right)^{1/2} \quad (0 \Leftrightarrow m_3 = m_8)$$

$$m_T = m_T^0 \left(\frac{m_3}{m_8} \right)^{5/2}$$

$$m_8 \cong M_{GUT}^0{}^2 / m_{pe} \Rightarrow M_{GUT} \cong 10 M_{GUT}^0$$

$$m_T \cong 30 m_T$$

$$d=5 \propto m_T^2 \cong 1000!$$

Bachas, Fabre,
Yanagida
(1986)

Bayc, Friolanes Perez,
G.S. (2002)

GIVE UP ON SUSY:

ORDINARY $SO(10)$

UNIFICATION: $M_{\text{SUT}(2)_R} \approx 10^{10} - 10^{13} \text{ GeV}$

$$\bullet \quad 16_H, 45_H \Rightarrow M_R = \frac{M_{\text{SUT}(2)_R}^2}{M_{\text{pl}}} \approx 10 - 10^7 \text{ GeV}$$

NOT GOOD

$$\bullet \quad \overline{126}_H, 210_H: M_R = 10^{10} - 10^{13} \text{ GeV}$$

GOOD

AF OK

In this case no problem
with large representations

of fields not an issue -

calculability

Bacella et al

(2000-2004)

Bay, Melzer, Visseri,

G.S.

(in preparation)